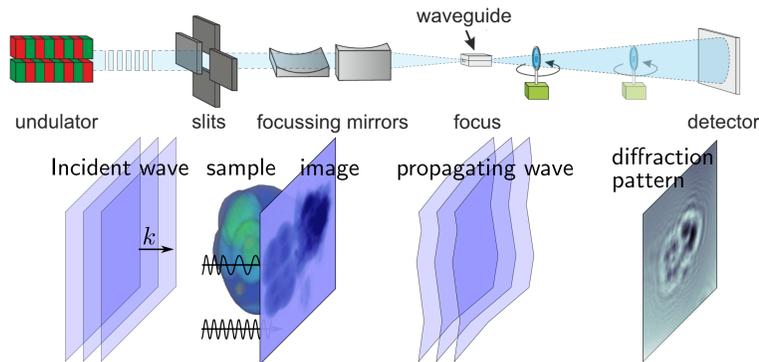


1. Introduction

Motivation: Imaging with X-rays permits nano-scale resolution owing to small wavelengths $\lesssim 1$ nm but suffers from weak absorption for microscopic samples.

↪ **Need for refraction-based imaging techniques** ↪ *phase contrast*

Imaging setup: Specimen is illuminated by coherent X-rays (e.g. from synchrotron [5]), resulting diffraction patterns are recorded downstream at propagation distance d :



Forward model: $F_{\text{PCI}} : n = \int_{-L}^0 (\delta - i\beta) dz \mapsto I_{\text{PCI}} = \left| \underbrace{\mathcal{D}}_{\text{Fresnel propagator}} \underbrace{(\exp(-ikn))}_{\text{exit wave field}} \right|^2$
 unknown image \mapsto measurable intensity data

► Model based on paraxial Helmholtz eq. + ray optics in sample \rightarrow valid for X-rays [1]

► Images yield 2D-projections of the object's complex refractive index $n = 1 - \delta + i\beta$

Phase contrast tomography: Image the refractive index in 3D by recording diffraction patterns under different incident angles θ (model: Radon transform \mathcal{R})

$$F_{\text{PCT}} : \tilde{n} = (\delta - i\beta) \mapsto \{I_{\text{PCT},\theta}\} = |\mathcal{D}(\exp(-ik\mathcal{R}(\tilde{n})))|^2$$

Inverse Problem: Reconstruct the 2D- or 3D-image n or \tilde{n} from observed noisy data $I^\varepsilon = I_* + \varepsilon$ and available *a priori* knowledge, i.e. invert the forward map F_* .

2. Phase Retrieval - an Unstable Procedure?

Phase retrieval: Image reconstruction requires inversion of $|\cdot|^2$, i.e. *phase recovery*, due to the physical restriction to measuring wave intensities

Is the image uniquely determined by the data? Yes! [3]

↪ **Is the reconstruction robust to measurement errors? ...**

Weak object approximation: Linearize F_{PCI} in $h := -ikn = -\mu - i\phi$ (valid if $\phi, \mu \ll 1$)

$$F_{\text{PCI}}(n) - 1 \approx \underbrace{\mathcal{D}(h)}_{\text{image}} + \underbrace{\mathcal{D}^{-1}(\bar{h})}_{\text{twin image}} =: T_{\text{PCI}}h \quad (\phi: \text{phase shifts}, \mu: \text{absorption})$$

Stability problem: Is there a constant $C_A > 0$ such that for all images $h_1, h_2 \in A \subset L^2(\mathbb{R}^2)$

$$C_A \frac{\|h_1 - h_2\|}{\text{object perturbation}} \leq \frac{\|T_{\text{PCI}}(h_1 - h_2)\|}{\text{increment in the data}} ? \Leftrightarrow \text{reconstruction error} \leq C_A^{-1} \|\varepsilon\|$$

Common a priori constraints – general instability:

► *Non-absorbing object:* $\mu = 0$ (biological tissue)

► *Single-material:* $\mu = \tan(\varphi_0)\phi$ (more general)

$$\Rightarrow T_{\text{PCI}}h \propto 2\mathcal{F}^{-1} \left(\underbrace{\sin\left(\frac{\xi^2}{4\pi N_F} + \varphi_0\right)}_{\text{contrast transfer fct. (CTF)}} \underbrace{\mathcal{F}(h)}_{\text{Fourier transform}} \right)$$

► Arbitrarily low contrast at CTF-zeros \rightsquigarrow bad SNR

► Unstable reconstruction without further constraints

↪ **How to choose the admissible objects A to ensure stability?**

References

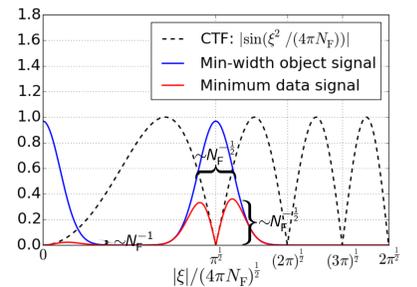
- [1] Jonas P and Louis A. Phase contrast tomography using holographic measurements. *Inverse Problems*, 20(1):75, 2004.
- [2] Maretzke S. Regularized Newton methods for simultaneous Radon inversion and phase retrieval in phase contrast tomography. *arXiv preprint*, page arXiv:1502.05073, 2015.
- [3] Maretzke S. A uniqueness result for propagation-based phase contrast imaging from a single measurement. *Inverse Problems*, 31:065003, 2015.
- [4] Maretzke S, Bartels M, Krenkel M, Salditt T, and Hohage T. Regularized Newton methods for X-ray phase contrast and general imaging problems. *Optics Express*, 2016. (accepted).
- [5] Salditt T, Osterhoff M, Krenkel M, Wilke R N, Priebe M, Bartels M, Kalbfleisch S, and Sprung M. Compound focusing mirror and x-ray waveguide optics for coherent imaging and nano-diffraction. *Journal of synchrotron radiation*, 22(4):867–878, 2015.
- [6] Slepian D and Sonnenblick E. Eigenvalues associated with prolate spheroidal wave functions of zero order. *Bell System Technical Journal*, 44(8):1745–1759, 1965.

3. Stability Estimates for Single-Material Objects

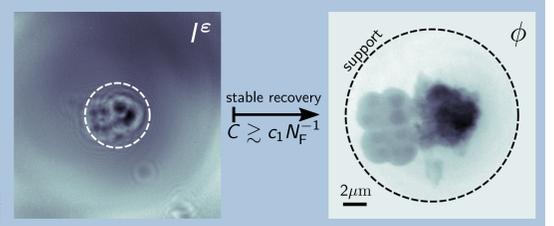
Support constraint: Object h is contained in a known subdomain $\Omega \subset \mathbb{R}^2$ of the field of view.

Uncertainty principle: Restriction $\text{supp}(h) \subset \Omega$ in real-space induces a *minimum* lengthscale in Fourier space:

$$\underbrace{\sigma_{|h|^2}}_{\leq \text{diam}(\Omega)} \cdot \sigma_{|\mathcal{F}(h)|^2} \gtrsim 1$$



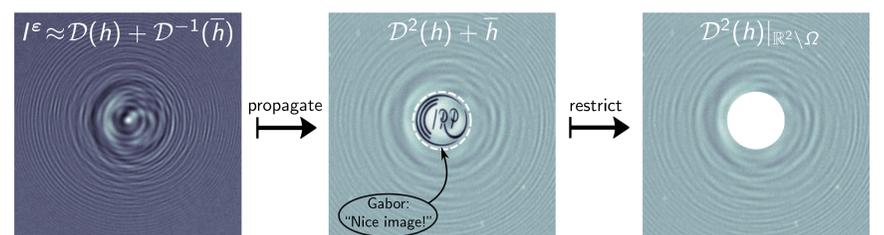
Theorem 1: Let N_F the Fresnel number of the support diameter $\text{diam}(\Omega)$. Then the recovery of *single-material objects* $h \propto \phi \in L^2(\Omega)$ is stable with $C_{\text{single}}(N_F) \geq \max \{ \min \{ C_1, c_1 N_F^{-1} \}, \min \{ C_2 \varphi_0^2, c_2 N_F^{-\frac{1}{2}} \} \}$



4. Stability Estimates for General Objects

↪ **Can we simultaneously reconstruct phase shifts ϕ and absorption μ ?**

Inverse Gabor holography: Propagate data and remove twin-image when in-focus



$$\|T_{\text{PCI}}h\|_{L^2} = \|\mathcal{D}T_{\text{PCI}}h\|_{L^2} = \|\mathcal{D}^2(h) + \bar{h}\|_{L^2} \geq \|\mathcal{D}^2(h)\|_{\mathbb{R}^m \setminus \Omega} = \|\mathcal{F}(m_F \cdot h)\|_{\mathbb{R}^2 \setminus (\pi N_F \Omega)} \|L^2$$

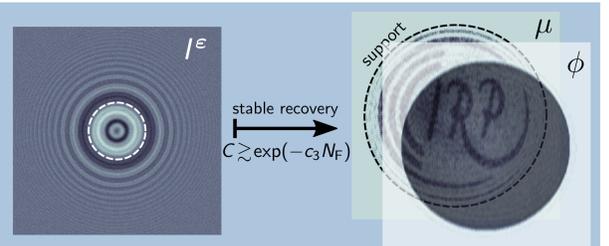
$$\Rightarrow C \geq \inf_{h \in L^2(\Omega), \|h\|=1} \|\mathcal{F}(h)\|_{\mathbb{R}^2 \setminus (\pi N_F \Omega)} \quad (N_F: \text{Fresnel number of } \Omega)$$

► Image recovery *at least as stable as* inverting $\mathcal{F}_F : L^2(\Omega) \rightarrow L^2(\mathbb{R}^2)$; $h \mapsto \mathcal{F}(h)|_{\mathbb{R}^2 \setminus (\pi N_F \Omega)}$

► Singular values of incomplete Fourier transform \mathcal{F}_F satisfy $\gtrsim \exp(-cN_F)$ [6] \rightsquigarrow *stable*

► Simultaneous recovery of ϕ and μ is feasible for deeply holographic measurements [4]

Theorem 2: Let N_F the Fresnel number of the support diameter $\text{diam}(\Omega)$. Then the recovery of *general objects* $h = -\mu - i\phi \in L^2(\Omega)$ is stable with $C_{\text{gen}}(N_F) \geq C_3 \exp(-c_3 N_F)$



5. Newton-Kaczmarz Methods for Phase Contrast Tomography

Standard method in phase contrast tomography:

- 1 *Independent* phase retrieval for all incident angles: $I_{\text{PCT}} = \{I_{\text{PCT},\theta}\} \xrightarrow{F_{\text{PCT}}^{-1}} \{n_\theta\} = \mathcal{R}(\delta - i\beta)$
- 2 *Subsequent* tomographic reconstruction by Radon inversion: $\mathcal{R}(\delta - i\beta) \mapsto \delta - i\beta$

Simultaneous approach: *All-at-once* inversion of F_{PCT} by regularized Newton methods

- *Stabilizing:* exploitation of tomographic correlations in phase reconstruction [2]
- *Efficient:* Process small subsets of incident angles per iteration \rightarrow *Newton-Kaczmarz* [4]
- *Flexible:* May account for non-ideal illumination, object motions, *misalignment*, etc.

$$\tilde{n}_{k+1} = \underset{\tilde{n}}{\text{argmin}} \left\| \underbrace{\mathcal{P}_k(F_{\text{PCT}}(\tilde{n}_k) + F'_{\text{PCT}}[\tilde{n}_k](\tilde{n} - \tilde{n}_k) - I^\varepsilon)}_{\text{restriction to small wedges of incident angles}} \right\|_{L^2}^2 + \underbrace{\alpha_k \|\tilde{n} - \tilde{n}_0\|_{X_k}^2}_{\text{approximate } L^p, \text{ TV, positivity, ...}} + \underbrace{\alpha_0 \|\tilde{n} - \tilde{n}_k\|_{L^2}^2}_{\text{strong regularization } \Rightarrow \text{well-conditioned}}$$

↪ **Efficient + accurate combination of phase retrieval and tomography**

